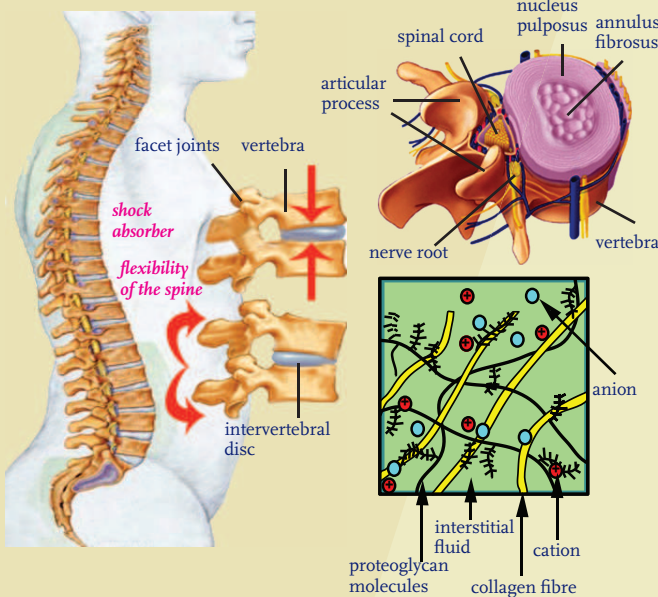


Introduction

The **swelling** and **shrinking** of cartilaginous tissues (CT) (like **Intervertebral Disks**) can be modelled by a four-component mixture theory in which a deformable and charged porous medium is saturated with a fluid with dissolved ions. This theory results in a coupled system of **non-linear parabolic** differential equations together with an algebraic constraint for **electro-neutrality**.

Four-Component Modeling of CT

- The two-component (Biot classical model) mixture theory is not able to describe the swelling and shrinking, that is caused by **chemical** and/or **electrical loads**.
- Therefore the theory is extended to **four-component** mixture theory.
- We assume **incompressible** and **linear elastic solid** saturated with **incompressible** and **Newtonian viscous** fluid. Also we neglect the inertial effects (quasistatic case) and the body forces.



Balance Equations

(degrees of freedom, material parameters)

$$\begin{cases} -\nabla \cdot (2\mu_s \mathbf{E}(\mathbf{u}) + \lambda_s \nabla \cdot \mathbf{u}) + \nabla p = \mathbf{0}, & \text{force balance} \\ \frac{\partial \nabla \cdot \mathbf{u}}{\partial t} + \nabla \cdot \mathbf{q}^l = 0, & \text{volume balance} \\ \frac{\partial c^\beta}{\partial t} (\nabla \cdot \mathbf{u} + \varphi_0) + \nabla \cdot (\mathbf{q}^\beta + c^\beta \mathbf{q}^l) = 0, \beta = +, -. & \text{ion balances} \end{cases}$$

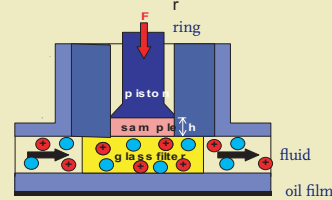
Constitutive Equations

$$\begin{cases} \mathbf{q}^l = -K(\nabla \mu^l + c^+ \nabla \mu^+ + c^- \nabla \mu^-), & \text{Darcy's law} \\ \mathbf{q}^\beta = -\frac{D^\beta c^\beta \varphi}{RT} \nabla \mu^\beta, \beta = +, -. & \text{Fick's law} \end{cases}$$

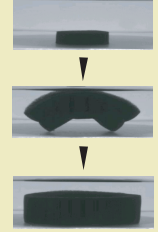
Secondary Equations

$$\begin{cases} \varphi = 1 - (1 - \varphi_0)(1 - \nabla \cdot \mathbf{u}), & c^{fc} = c_0^{fc} (1 - \nabla \cdot \mathbf{u} / \varphi_0), \\ c^\beta = \frac{1}{2z^\beta} c^{fc} + \frac{1}{2} \sqrt{(c^{fc})^2 + 4c^2 \exp\left(\frac{\mu^+ - \mu_0^+ + \mu^- - \mu_0^-}{RT}\right)}, \\ p = \mu^f - \mu_0^f + RT(\Gamma^+ c^+ + \Gamma^- c^-), \\ \xi = \frac{1}{z^\beta F} (\mu^\beta - \mu_0^\beta - RT \ln \frac{c^\beta}{c}), \quad \beta = +, -. \end{cases}$$

Experiments and Boundary Conditions



Experimental Set-Up



Experiment on hydrogel disc J.M. Huyghe 1999

Mixed Finite Element

- For the sake of **local mass conservation**, equations are discretised in space using mixed finite element.
- We choose $\mathbf{u} \in \mathcal{P}_1(\Omega)^2$, \mathbf{q}^l , $\mathbf{q}_{tot}^\beta := \mathbf{q}^\beta + c^\beta \mathbf{q}^l \in \mathcal{RT}_0(\Omega)$ and $\mu^l, \mu^\beta \in \mathcal{P}_0(\Omega)$.
- Implicit time integration results into a non-linear set of equations.
- Since different scales are apparent in the weak formulation, then we replace the reduced matrices by the scaled ones.
- Substitution in the discrete variational formulation gives:

$$\mathfrak{A}(\varphi, c^+, c^-) \frac{dy}{dt} + \mathfrak{B}(\varphi, c^+, c^-) y = \mathfrak{F}(\varphi_h, c_h^+, c_h^-) + \frac{d\mathfrak{G}}{dt}(\varphi_h, c_h^+, c_h^-),$$

$$y = [\bar{u}, \bar{q}^l, \bar{q}_{tot}^+, \bar{q}_{tot}^-, \bar{\mu}^l, \bar{\mu}^+, \bar{\mu}^-]^T,$$

$$\mathfrak{A}_{ij} = \begin{cases} \mathbf{B}^T, & i = 5, j = 1, \\ 0, & i \neq 5, j \neq 1, \end{cases}$$

$$\mathfrak{B} = \begin{pmatrix} \mathbf{A} & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{B} & \mathbf{0} & \mathbf{0} \\ 0 & \mathbf{C}^{ll}(\varphi_h/c_h^\beta) & \mathbf{C}^{l+}(\varphi_h) & \mathbf{C}^{l-}(\varphi_h) & \mathbf{D} & \mathbf{0} & \mathbf{0} \\ 0 & \mathbf{C}^{l+}(\varphi_h) & \mathbf{C}^{++}(\varphi_h, c_h^+) & \mathbf{0} & \mathbf{0} & \mathbf{D} & \mathbf{0} \\ 0 & \mathbf{C}^{l-}(\varphi_h) & \mathbf{0} & \mathbf{C}^{--}(\varphi_h, c_h^-) & \mathbf{0} & \mathbf{0} & \mathbf{D} \\ 0 & \mathbf{D}^T & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ 0 & \mathbf{0} & \mathbf{D}^T & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ 0 & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \end{pmatrix} \oplus_{3K \times 3K}$$

- The upper-left block is symmetric positive definite.
- Given φ_h and c_h^β , the solution for the linear saddle point problem exists and it is unique.

Analytical Solution

- To verify the numerical solutions we derived a set of analytical solutions for the one-dimensional problem.
- The idea behind is to derive a coupled system of diffusion equations.

$$\frac{\partial}{\partial t} \begin{pmatrix} \mu^l \\ \mu^+ \\ \mu^- \end{pmatrix} = \mathbf{E} \mathbf{P} \frac{\partial^2}{\partial x^2} \begin{pmatrix} \mu^l \\ \mu^+ \\ \mu^- \end{pmatrix}$$

- \mathbf{E} and \mathbf{P} are both symmetric positive definite and related to the Hessian of Helmholtz free energy and the diffusion matrix, respectively.

References

- Malakpoor et al., An analytical solution of incompressible charged porous media , ZAMM. Z. Angew. Math. Mech, 2006.
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