# Problem Companion <br> Tensor Calculus and Differential Geometry 

2WAHO



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Cover illustration: papyrus fragment from Euclid's Elements of Geometry, Book II [8].

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## Preface

This problem companion belongs to the course notes "Tensor Calculus and Differential Geometry" (course code 2WAH0) by Luc Florack. Problems are concisely formulated. Einstein summation convention applies to all problems, unless stated otherwise. Please refer to the course notes for further details.

## 1. Prerequisites from Linear Algebra

1. Let $V$ be a real vector space, and $u \in V$. Show that $(-1) \cdot u=-u$, and $0 \cdot u=o$.
2. Show that $\mathscr{L}(V, W)$ is a vector space.

In analogy with the determinant of a matrix $A$, the so-called permanent is defined as

$$
\operatorname{perm} A=\sum_{j_{1}, \ldots, j_{n}=1}^{n}\left|\left[j_{1}, \ldots, j_{n}\right]\right| A_{1 j_{1}} \ldots A_{n j_{n}}=\frac{1}{n!} \sum_{\substack{i_{1}, \ldots, i_{n}=1 \\ j_{1}, \ldots, j_{n}=1}}^{n}\left|\left[i_{1}, \ldots, i_{n}\right]\left[j_{1}, \ldots, j_{n}\right]\right| A_{i_{1} j_{1}} \ldots A_{i_{n} j_{n}}
$$

Note that the nontrivial factors among the weights $\left|\left[i_{1}, \ldots, i_{n}\right]\right|$ and $\left|\left[j_{1}, \ldots, j_{n}\right]\right|$ are invariably +1 .
3. Can we omit the factors $\left|\left[i_{1}, \ldots, i_{n}\right]\right|$ and $\left|\left[j_{1}, \ldots, j_{n}\right]\right|$ in this definition of perm $A$ ?
4. How many multiplications and additions/subtractions do you need in the numerical computation of $\operatorname{det} A$ and $\operatorname{perm} A$ for a generic $n \times n$ matrix $A$ ?
5. Given a generic $n \times n$ matrix $A$. Argue why its cofactor and adjugate matrices $\tilde{A}$, respectively $\tilde{A}^{\mathrm{T}}$, always exist, unlike its inverse $A^{-1}$. What is the condition for $A^{-1}$ to be well-defined?
6. Let

$$
A=\left(\begin{array}{ll}
1 & 3 \\
2 & 7
\end{array}\right) \quad \text { and } \quad B=\left(\begin{array}{ccc}
1 & 2 & 1 \\
6 & 3 & 8 \\
8 & 7 & 10
\end{array}\right)
$$

Compute, by hand, $\operatorname{det} A$ and $\operatorname{det} B$. Likewise for perm $A$ and perm $B$.
7. Cf. previous problem. Compute the following expression involving the standard 3-dimensional inner and outer product of vectors (cf. the columns of $B$ ), and provide a geometrical interpretation of your result for det $B$ :

$$
\left[\left(\begin{array}{l}
1 \\
6 \\
8
\end{array}\right) \times\left(\begin{array}{l}
2 \\
3 \\
7
\end{array}\right)\right] \cdot\left(\begin{array}{c}
1 \\
8 \\
10
\end{array}\right)
$$

8. Cf. previous problem. Compute the cofactor and adjugate matrices $\tilde{A}, \tilde{B}, \tilde{A}^{\mathrm{T}}$, and $\tilde{B}^{\mathrm{T}}$.
9. Cf. previous problem. Compute the inverse matrices $A^{-1}$ and $B^{-1}$, if these exist.
10. Cf. previous problem. Compute $A \tilde{A}^{\mathrm{T}}$ and $B \tilde{B}^{\mathrm{T}}$.
11. Given $\operatorname{det} A$, $\operatorname{det} B$ for general $n \times n$ matrices $A, B$. What is $\operatorname{det} \tilde{A}, \operatorname{det} A^{\mathrm{T}}, \operatorname{det}(\lambda A), \operatorname{det}(A B)$, and $\operatorname{det} A^{k}$ for $\lambda \in \mathbb{R}, k \in \mathbb{Z}$ ?
12. Consider the collection $A_{i_{1} \ldots i_{n}}$ for all $i_{1}, \ldots, i_{n}=1, \ldots, n$. Suppose $A_{i_{1} \ldots i_{k} \ldots i_{\ell} \ldots i_{n}}=-A_{i_{1} \ldots i_{\ell} \ldots i_{k} \ldots i_{n}}$ for any $1 \leq k<\ell \leq n$ (complete antisymmetry). Show that $A_{i_{1} \ldots i_{n}} \propto\left[i_{1} \ldots i_{n}\right]$.
13. Prove the following identities for the completely antisymmetric symbol in $n=3$ :

$$
\begin{aligned}
{[i, j, k][\ell, m, n] } & =\delta_{i \ell} \delta_{j m} \delta_{k n}+\delta_{i m} \delta_{j n} \delta_{k \ell}+\delta_{i n} \delta_{j \ell} \delta_{k m}-\delta_{i m} \delta_{j \ell} \delta_{k n}-\delta_{i \ell} \delta_{j n} \delta_{k m}-\delta_{i n} \delta_{j m} \delta_{k \ell} \\
\sum_{i=1}^{3}[i, j, k][i, m, n] & =\delta_{j m} \delta_{k n}-\delta_{j n} \delta_{k m} \\
\sum_{i, j=1}^{3}[i, j, k][i, j, n] & =2 \delta_{k n} \\
\sum_{i, j, k=1}^{3}[i, j, k][i, j, k] & =6
\end{aligned}
$$

$\star$ 14. Compute the Gaussian integral $\gamma(A)=\int_{\mathbb{R}^{n}} \exp \left(-x^{i} A_{i j} x^{j}\right) d x$ for a symmetric positive definite $n \times n$ matrix $A$. Hint: There exists a rotation matrix $R$ such that $R^{\mathrm{T}} A R=\Delta$, in which $\Delta=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$, with all $\lambda_{p}>0, p=1, \ldots, n$.
$\star$ 15. Compute the extended Gaussian integral $\gamma(A, s)=\int_{\mathbb{R}^{n}} \exp \left(-x^{i} A_{i j} x^{j}+s_{k} x^{k}\right) d x$ for a symmetric positive definite $n \times n$ matrix $A$ and arbitrary "source" $s \in \mathbb{R}^{n}$.
$\star$ 16. Cf. previous problem. Consider the following integral: $\gamma^{i_{1} \ldots i_{p}}(A)=\int_{\mathbb{R}^{n}} x^{i_{1}} \ldots x^{i_{p}} \exp \left(-x^{i} A_{i j} x^{j}\right) d x$. Express $\gamma^{i_{1} \ldots i_{p}}(A)$ in terms of $\gamma(A, s)$. (You don't need to compute $\gamma^{i_{1} \ldots i_{p}}(A)$ explicitly.)

## 2. Tensor Calculus

17. Expand in $n=2$ and $n=3$ dimensions, respectively: $X_{i} Y^{i j} X_{j}$. Argue why we may assume, without loss of generality, that $Y^{i j}=Y^{j i}$ in this expression ("automatic" symmetry).
18. Expand in $n=2$ and $n=3$ dimensions, respectively: $X_{i j} Y^{i j}$. Are we allowed to assume that $X_{i j}$ or $Y^{i j}$ is symmetric in this expression?
19. Show that $\delta_{i}^{i}=n$ and $\delta_{k}^{i} \delta_{j}^{k}=\delta_{j}^{i}$.
20. Show that $\delta_{j}^{i} X_{i}^{j}=X_{i}^{i}$.
21. Suppose $x^{i}=A_{j}^{i} y^{j}$ and $y^{i}=B_{j}^{i} x^{j}$ for all $y \in \mathbb{R}^{n}$. Prove that $A_{k}^{i} B_{j}^{k}=\delta_{j}^{i}$.
22. Consider a Cartesian basis $\left\{\partial_{1} \equiv \partial_{x}, \partial_{2} \equiv \partial_{y}\right\}$ spanning a 2-dimensional plane. (Here $\partial_{i}$ is shorthand for $\partial / \partial x^{i}$ if $x^{i}$ denotes the $i$-th Cartesian coordinate; in $\mathbb{R}^{2}$ we identify $x^{1} \equiv x, x^{2} \equiv y$.) Let $\bar{x}^{i}$ denote the $i$-th polar coordinate, with $\bar{x}^{1} \equiv r$ (radial distance) and $\bar{x}^{2} \equiv \phi$ (polar angle), such that $x=r \cos \phi, y=r \sin \phi$. Assume $\bar{\partial}_{j}=A_{j}^{i} \partial_{i}$ (with $\bar{\partial}_{1} \equiv \partial_{r}, \bar{\partial}_{2} \equiv \partial_{\phi}$ ). Argue why this assumption holds, and compute the matrix $A$.
23. Cf. previous problem, but now for the relation between Cartesian and spherical coordinates in $\mathbb{R}^{3}$, defined by $x=r \sin \theta \cos \phi, y=r \sin \theta \sin \phi, z=r \cos \theta$, with radial distance $\bar{x}^{1} \equiv r$, polar angle $\bar{x}^{2} \equiv \theta$ and azimuthal angle $\bar{x}^{3} \equiv \phi$, and, correspondingly, $\bar{\partial}_{1} \equiv \partial_{r}, \bar{\partial}_{2} \equiv \partial_{\theta}, \bar{\partial}_{3} \equiv \partial_{\phi}$.
24. Cf. previous two problems. Consider the dual Cartesian bases $\left\{d x^{1} \equiv d x, d x^{2} \equiv d y\right\}$ in two dimensions, respectively $\left\{d x^{1} \equiv d x, d x^{2} \equiv d y,, d x^{3} \equiv d z\right\}$ in three dimensions. Assume $d x^{i}=C_{j}^{i} d \bar{x}^{j}$, with dual bases $\left\{d \bar{x}^{1} \equiv d r, d \bar{x}^{2} \equiv d \phi\right\}$, respectively $\left\{d \bar{x}^{1} \equiv d r, d \bar{x}^{2} \equiv d \theta, d \bar{x}^{3} \equiv d \phi\right\}$. Show that $C=A$, i.e. the same matrix as computed in the previous two problems. (Notice the difference!)
25. Let $\mathbf{x}=x^{i} \mathbf{e}_{i} \in V$. Show that the map $\hat{\mathbf{a}}: V \rightarrow \mathbb{R}$ defined by $\hat{\mathbf{a}}(\mathbf{x})=a_{i} x^{i}$ is a linear operator.
26. Cf. previous problem. A linear operator of this type is known as a covector, notation $\hat{\mathbf{a}} \in V^{*} \equiv \mathscr{L}(V, \mathbb{R})$. Explain why a covector â (respectively covector space $V^{*}$ ) is formally a vector (respectively vector space).
27. Consider $\hat{\mathbf{a}} \in V^{*}$ with prototype $\hat{\mathbf{a}}: V \rightarrow \mathbb{R}: \mathbf{x} \mapsto \hat{\mathbf{a}}(\mathbf{x})=a_{i} x^{i}$. Argue why this naturally provides an alternative interpretation of $\mathbf{x} \in V$ as an element of $V^{* *}=\left(V^{*}\right)^{*} \sim V$.
28. Suppose $\langle\hat{\omega}, \mathbf{v}\rangle=C_{j}^{i} \omega_{i} v^{j}=\bar{C}_{j}{ }_{j} \bar{\omega}_{i} \bar{v}^{j}$ relative to dual bases $\left\{\mathbf{e}_{i}, \hat{\mathbf{e}}^{i}\right\}$ (middle term), respectively $\left\{\mathbf{f}_{i}, \hat{\mathbf{f}}^{i}\right\}$ (right hand side), in which $C_{j}^{i}$ and $\bar{C}_{j}^{i}$ are coefficients ("holors") to be determined. Show that the holor is basis independent, i.e. $C_{j}^{i}=\bar{C}_{j}^{i}$, and compute its components explicitly.
29. Suppose $\left\{\mathbf{e}_{i}\right\}$ and $\left\{\hat{\mathbf{e}}^{i}\right\}$ are dual bases of $V$, respectively $V^{*}$, and $\mathbf{e}_{i}=A_{i}^{j} \mathbf{f}_{j}, \hat{\mathbf{e}}^{i}=B_{j}^{i} \hat{\mathbf{f}}^{j}$ for some (a priori
unrelated) transformation matrices $A$ and $B$. Show that if $\left\{\mathbf{f}_{j}, \hat{\mathbf{f}}^{j}\right\}$ constitute dual bases then $B_{k}^{j} A_{i}^{k}=\delta_{i}^{j}$.
$\star$ 30. Show that the following "length functional" $\mathscr{L}(\gamma)$ for a parameterized curve $\gamma:\left[T_{-}, T_{+}\right] \rightarrow \mathbb{R}^{n}: t \mapsto \gamma(t)$ with fixed end points $X_{ \pm}=\gamma\left(T_{ \pm}\right)$is independent of the parametrization:

$$
\mathscr{L}(\gamma)=\int_{T_{-}}^{T_{+}} \sqrt{g_{i j}(\gamma(t)) \dot{\gamma}^{i}(t) \dot{\gamma}^{j}(t)} d t
$$

Here $\dot{\gamma}(t)=\dot{\gamma}^{i}(t) \mathbf{e}_{i}$ denotes the derivative of the curve, expanded relative to a fiducial basis $\left\{\mathbf{e}_{i}\right\}$, and $\gamma_{i j}(x)$ are the components of the inner product defined at base point $x \in \mathbb{R}^{n}$. To this end, consider a reparametrization of the form $s=s(t)$, with $\dot{s}(t)>0$, say.
31. Cf. previous problem. The parameter $s$ is called an affine parameter if, along the entire parameterized curve $\xi:\left[S_{-}, S_{+}\right] \rightarrow \mathbb{R}^{n}: s \mapsto \xi(s)$, we have $\|\dot{\xi}(s)\|=1$ ("unit speed parameterization"), or $(\dot{\xi}(s) \mid \dot{\xi}(s))=1$, in which the l.h.s. pertains to the inner product at point $\xi(s) \in \mathbb{R}^{n}$. In other words, $g_{i j}(\xi(s)) \dot{\xi}^{i}(s) \dot{\xi}^{j}(s)=1$. Show that this can always be realized through suitable reparameterization starting from an arbitrarily parameterized curve $\gamma:\left[T_{-}, T_{+}\right] \rightarrow \mathbb{R}^{n}: t \mapsto \gamma(t)$.
32. The figure below shows a pictorial representation of vectors $(\mathbf{v}, \mathbf{w} \in V)$ and covectors $\left(\boldsymbol{\omega} \in V^{*}\right)$ in terms of graphical primitives. In this picture, a vector is denoted by a directed arrow, and a covector by an equally spaced set of level lines along a directed normal (i.e. "phase" increases in the direction of the directed normal, attaining consecutive integer values on the level lines drawn in the figure). Give a graphical interpretation of the contraction $\langle\boldsymbol{\omega}, \mathbf{v}\rangle \in \mathbb{R}$, and estimate from the figure the values of $\langle\boldsymbol{\omega}, \mathbf{v}\rangle,\langle\boldsymbol{\omega}, \mathbf{w}\rangle$, and $\langle\boldsymbol{\omega}, \mathbf{v}+\mathbf{w}\rangle$. Are these values consistent with the linearity of $\langle\boldsymbol{\omega}, \cdot\rangle$ ?

33. An inner product on $V$ induces an inner product on $V^{*}$. Recall that for $\mathbf{x}=x^{i} \mathbf{e}_{i} \in V$ and $\mathbf{y}=y^{j} \mathbf{e}_{j} \in V$ we have $(\mathbf{x} \mid \mathbf{y})=g_{i j} x^{i} y^{j}$. Let $\hat{\mathbf{x}}=x_{i} \hat{\mathbf{e}}^{i} \in V^{*}, \hat{\mathbf{y}}=y_{j} \hat{\mathbf{e}}^{j} \in V^{*}$, with $\left\langle\hat{\mathbf{e}}^{i}, \mathbf{e}_{j}\right\rangle=\delta_{j}^{i}$. Define $(\hat{\mathbf{x}} \mid \hat{\mathbf{y}})_{*}=(b \hat{\mathbf{x}} \mid b \hat{\mathbf{y}})$. Show that $(\hat{\mathbf{x}} \mid \hat{\mathbf{y}})_{*}=g^{i j} x_{i} y_{j}$, in which $g^{i k} g_{k j}=\delta_{j}^{i}$.
34. Let $\hat{\mathbf{u}}, \hat{\mathbf{v}} \in V^{*}$. Prove equivalence of nilpotency and antisymmetry of the wedge product, i.e. $\hat{\mathbf{u}} \wedge \hat{\mathbf{u}}=0$ iff $\hat{\mathbf{u}} \wedge \hat{\mathbf{v}}=-\hat{\mathbf{v}} \wedge \hat{\mathbf{u}}$.
35. Let $\pi, \pi^{\prime}:\{1, \ldots, p\} \rightarrow\{1, \ldots, p\}: k \mapsto \pi(k)$ be two bijections. By abuse of notation these can be identified with permutations on any symbolic set $\Omega_{p}=\left\{\left(a_{1}, \ldots, a_{p}\right)\right\}$ consisting of $p$-tuples of labeled symbols: $\pi: \Omega_{p} \rightarrow \Omega_{p}:\left(a_{1}, \ldots, a_{p}\right) \mapsto\left(a_{\pi(1)}, \ldots, a_{\pi(p)}\right)$, and likewise for $\pi^{\prime}$. Let $\pi(1, \ldots, p)=(\pi(1), \ldots, \pi(p))$, respectively $\pi^{\prime}(1, \ldots, p)=\left(\pi^{\prime}(1), \ldots, \pi^{\prime}(p)\right)$, argue that $\pi^{\prime \prime}=\pi^{\prime} \pi$, i.e. the (right-to-left) concatenation of $\pi^{\prime}$ and $\pi$, defines another permutation.
36. Argue that the concatenation $\pi^{\prime \prime}=\pi^{\prime} \pi$ of two permutations $\pi^{\prime}$ and $\pi$ satisfies $\operatorname{sgn} \pi^{\prime \prime}=\operatorname{sgn} \pi^{\prime} \operatorname{sgn} \pi$.
37. Show that the symmetrisation map $\mathscr{S}: \mathbf{T}_{p}^{0}(V) \rightarrow \bigvee_{p}(V)$ is idempotent, i.e. $\mathscr{S} \circ \mathscr{S}=\mathscr{S}$.
38. Show that the antisymmetrisation $\operatorname{map} \mathscr{A}: \mathbf{T}_{p}^{0}(V) \rightarrow \bigwedge_{p}(V)$ is idempotent, i.e. $\mathscr{A} \circ \mathscr{A}=\mathscr{A}$.
39. Cf. the two previous problems. Show that $\mathscr{A} \circ \mathscr{S}=\mathscr{S} \circ \mathscr{A}=0 \in \mathscr{L}\left(\mathbf{T}_{p}^{0}(V), \mathbf{T}_{p}^{0}(V)\right)$, i.e. the null operator on $\mathbf{T}_{p}^{0}(V)$ for $p \geq 2$. What if $p=0,1$ ?
40. Let $t_{i_{1} \ldots i_{p}}$ be the holor of $\mathbf{T} \in \mathbf{T}_{p}^{0}(V)$. Show that the holor of $\mathscr{S}(\mathbf{T}) \in \bigwedge_{p}(V)$ is $t_{\left(i_{1} \ldots i_{p}\right)} \stackrel{\text { def }}{=} \frac{1}{p!} \sum_{\pi} t_{\pi\left(i_{1}\right) \ldots \pi\left(i_{p}\right)}$.
41. Cf. previous problem. Show that the holor of $\mathscr{A}(\mathbf{T}) \in \bigwedge_{p}(V)$ is $t_{\left[i_{1} \ldots i_{p}\right]} \stackrel{\text { def }}{=} \frac{1}{p!} \sum_{\pi} \operatorname{sgn} \pi t_{\pi\left(i_{1}\right) \ldots \pi\left(i_{p}\right)}$.
42. Write out the symmetrised holor $t_{(i j k)}$ in terms of $t_{i j k}$. Likewise for the antisymmetrised holor $t_{[i j k]}$.
43. Simplify the expressions of the previous problem as far as possible if $t_{i j k}$ is known to possess the partial symmetry property $t_{i j k}=t_{j i k}$.
44. Let $\mathscr{B}=\{d t, d x, d y, d z\}$ be a coordinate basis of the 4-dimensional spacetime covector space $V^{*}$ of $V \sim \mathbb{R}^{4}$. What is $\operatorname{dim} \bigwedge_{p}(V)$ for $p=0,1,2,3,4$ ? Provide explicit bases $\mathscr{B}_{p}$ of $\bigwedge_{p}(V)$ for $p=0,1,2,3,4$ induced by $\mathscr{B}$.
45. Using the definition $\left(\mathbf{v}_{1} \wedge \ldots \wedge \mathbf{v}_{k} \mid \mathbf{x}_{1} \wedge \ldots \wedge \mathbf{x}_{k}\right)=\operatorname{det}\left\langle\sharp \mathbf{v}_{i}, \mathbf{x}_{j}\right\rangle$, show that for any $\mathbf{v}, \mathbf{w} \in V$ (with $\operatorname{dim} V \geq 2),(\mathbf{v} \wedge \mathbf{w} \mid \mathbf{v} \wedge \mathbf{w}) \geq 0$, with equality iff $\mathbf{v} \wedge \mathbf{w}=0$. What if $\operatorname{dim} V=1$ ?

Hint: Recall the Schwartz inequality for inner products: $|(\mathbf{v} \mid \mathbf{w})| \leq \sqrt{(\mathbf{v} \mid \mathbf{v})(\mathbf{w} \mid \mathbf{w})}$.
46. Cf. previous problem. Show that $(\mathbf{v} \wedge \mathbf{w} \mid \mathbf{x} \wedge \mathbf{y})=(\mathbf{x} \wedge \mathbf{y} \mid \mathbf{v} \wedge \mathbf{w})$ for all $\mathbf{v}, \mathbf{w}, \mathbf{x}, \mathbf{y} \in V$.
47. Expand and simplify the symmetrized holor $g_{(i j} h_{k \ell)}$ in terms of the unsymmetrized holor, i.e. in terms of terms like $g_{i j} h_{k \ell}$ and similar ones with permuted index orderings, using the symmetry properties $g_{i j}=g_{j i}$ and $h_{i j}=h_{j i}$.
48. Show that for the holor of any covariant 2-tensor we have $T_{i j}=T_{(i j)}+T_{[i j]}$.
49. Cf. the previous problem. Show that, for the holor of a typical covariant 3-tensor, $T_{i j k} \neq T_{(i j k)}+T_{[i j k]}$.
50. Expand and simplify the symmetrized holor $g_{(i j} g_{k \ell)}$ in terms of the unsymmetrized holor, i.e. in terms of terms like $g_{i j} g_{k \ell}$ and similar ones with permuted index orderings, using the symmetry property $g_{i j}=g_{j i}$.
51. Consider the 1-form $d f=\partial_{i} f d x^{i}$ and define the gradient vector $\nabla f=b d f=\partial^{i} f \partial_{i f}$ relative to the corresponding dual bases $\left\{d x^{i}\right\}$ and $\left\{\partial_{i}\right\}$. Here $\partial_{i} f$ is shorthand for the partial derivative $\frac{\partial f}{\partial x^{i}}$, whereas $\partial^{i} f$ is a symbolic notation for the contravariant coefficient of $\nabla f$ relative to $\left\{\partial_{i}\right\}$ (all evaluated at a fiducial point). Show that $\partial^{i} f=g^{i j} \partial_{j} f$.
52. Cf. previous problem.

- Compute the matrix entries $g^{i j}$ explicitly for polar coordinates in $n=2$, i.e. $\bar{x}^{1} \equiv r$ and $\bar{x}^{2} \equiv \phi$, starting from a standard inner product in Cartesian coordinates $x^{1} \equiv x=r \cos \phi, x^{2}=y=r \sin \phi$. Relative to basis $\left\{\partial_{i}\right\} \widehat{=}\left\{\partial_{x}, \partial_{y}\right\}$, in which $\mathbf{v}=v^{i} \partial_{i}=v^{x} \partial_{x}+v^{y} \partial_{y}$, respectively $\mathbf{w}=w^{i} \partial_{i}=w^{x} \partial_{x}+w^{y} \partial_{y}$, the standard inner product takes the form $(\mathbf{v} \mid \mathbf{w})=\eta_{i j} v^{i} w^{j}$ with $\eta_{i j}=1$ if $i=j$ and 0 otherwise.
- Using your result, compute the components $\bar{\partial}^{1} f=\partial^{r} f$ and $\bar{\partial}^{2} f=\partial^{\phi} f$ of the gradient vector $\nabla f=$ $b d f=\bar{\partial}^{i} f \bar{\partial}_{i}$ relative to the polar basis $\left\{\bar{\partial}_{i}\right\} \widehat{=}\left\{\partial_{r}, \partial_{\phi}\right\}$.

53. Show that the Kronecker symbol $\delta_{j}^{i}$ is an invariant holor under the "tensor transformation law", and explain the attribute "invariant" in this context.
54. Cf. previous problem. Define $\boldsymbol{\epsilon}=\sqrt{g} \boldsymbol{\mu} \in \bigwedge_{n}(V)$, in which $V$ is an inner product space, $\boldsymbol{\mu}=d x^{1} \wedge \ldots \wedge d x^{n}$, and $g=\operatorname{det} g_{i j}$, the determinant of the covariant metric tensor. Consider its expansion relative to two bases, $\left\{d x^{i}\right\}$ respectively $\left\{d \bar{x}^{i}\right\}$, viz. $\boldsymbol{\epsilon}=\epsilon_{i_{1} \ldots i_{n}} d x^{i_{1}} \otimes \ldots \otimes d x^{i_{n}}=\bar{\epsilon}_{i_{1} \ldots i_{n}} d \bar{x}^{i_{1}} \otimes \ldots \otimes d \bar{x}^{i_{n}}$. Find the relation between the respective holors $\epsilon_{i_{1} \ldots i_{n}}$ and $\bar{\epsilon}_{i_{1} \ldots i_{n}}$. Is the holor of the $\boldsymbol{\epsilon}$-tensor invariant in the same sense as above?

## 3. Differential Geometry

55. Show that the definitions (i) $\nabla_{\partial_{j}} \partial_{i}=\Gamma_{i j}^{k} \partial_{k}$ and (ii) $\Gamma_{i j}^{k}=\left\langle d x^{k}, \nabla_{\partial_{j}} \partial_{i}\right\rangle$ are equivalent.
56. Show that if $\Gamma_{i j}^{k}=\Gamma_{j i}^{k}$ in $x$-coordinates, then $\bar{\Gamma}_{i j}^{k}=\bar{\Gamma}_{j i}^{k}$ in any coordinate basis induced by an arbitrary coordinate transformation $x=x(\bar{x})$.
57. Show that the holor $T_{i j}^{k}=\Gamma_{j i}^{k}-\Gamma_{i j}^{k}$ "transforms as a $\binom{1}{2}$-tensor".
58. Show that the holor $t_{\ell}=\Gamma_{k \ell}^{k}$ does not "transform as a covector", and derive its proper transformation law.
59. Consider 3-dimensional Euclidean space E furnished with Cartesian coordinates $(x, y, z) \in \mathbb{R}^{3}$. Assuming all Christoffel symbols $\Gamma_{i j}^{k}$ vanish relative to the induced Cartesian basis $\left(\partial_{x}, \partial_{y}, \partial_{z}\right)$, derive the corresponding symbols $\bar{\Gamma}_{i j}^{k}$ relative to a cylindrical coordinate system

$$
\left\{\begin{array}{l}
x=\rho \cos \phi \\
y=\rho \sin \phi \\
z=\zeta
\end{array}\right.
$$

60. Cf. the previous problem. Compute the components $g_{i j}$ of the metric tensor for Euclidean 3-space E in cylindrical coordinates $y^{i}$, assuming a Euclidean metric $g_{i j}(y) d y^{i} \otimes d y^{j}=\eta_{i j} d x^{i} \otimes d x^{j}$ in which the $x^{i}$ are Cartesian coordinates, and $\eta_{i j}=1$ if $i=j$ and $\eta_{i j}=0$ otherwise.
61. Cf. previous problems. Show that the Christoffel symbols $\bar{\Gamma}_{i j}^{k}$ derived above are in fact those corresponding to the Levi-Civita connection of Euclidean 3-space E (in cylindrical coordinates), and provide the geodesic equations for a geodesic curve $(\rho, \phi, \zeta)=(\rho(t), \phi(t), \zeta(t))$ in cylindrical coordinates. Finally, show that the solutions to these equations are indeed straight lines.

Hint: Compare the $y$-geodesic equations to the trivial $x$-geodesic equations $\ddot{x}=\ddot{y}=\ddot{z}=0$.

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